

# Formulating Quantum Theory as a Causally Neutral Theory of Bayesian Inference

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- 2 Spatial Quantum Conditional States
- 3 Temporal Quantum Conditional States
- 4 Preparations and Measurements
- 5 Quantum Bayes' rule
- 6 Bayesian Conditioning
- 7 Combining Expert Advice
- 8 Further results and open questions

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- The Church of the Larger Hilbert Space (J. Smolin)
  - Quantum theory is “about” a pure state vector of the universe that evolves unitarily.
  - Schrödinger, Everett, Zurek, . . .
- The Church of the Smaller Hilbert Space
  - Quantum theory is a noncommutative generalization of classical probability theory.
  - Heisenberg, von Neumann, . . .

**The Church of the Larger Hilbert Space** 







<http://www.facebook.com/group.php?gid=5658946617>

**The Church of The Smaller Hilbert Space** 

<http://www.facebook.com/group.php?gid=5965533115>







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- Classical probability theory does not care about causality
  - $P(Z, W, \dots)$
- Conventional quantum formalism does. . .

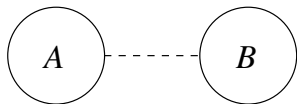


Figure: “Spacelike” correlations

$$\rho_{AB}$$

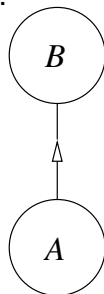


Figure: “Timelike” correlations

$$\rho_B = \mathcal{E}_{B|A}(\rho_A)$$

- Conventional Formalism: Hilbert spaces are attached to systems that persist in time.
  - States are a catalogue of probabilities for potential future measurement outcomes.
- Conditional States Formalism: Hilbert spaces are attached to systems at a specific time, or more generally to spacetime **regions**.
  - Always use a distinct label to distinguish input and output systems of a channel.
  - Always combine regions via the tensor product.
  - States are a catalogue of probabilities for any classical variables correlated with the region.

Table: Basic definitions

Classical Probability	Quantum Theory
Sample space $\Omega_Z = \{1, 2, \dots, d_Z\}$	Hilbert space $\mathcal{H}_A = \mathbb{C}^{d_A}$ $= \text{span}( 1\rangle,  2\rangle, \dots,  d_A\rangle)$
Probability distribution $P(Z = z) \geq 0$ $\sum_{z \in \Omega_Z} P(Z = z) = 1$	Quantum state $\rho_A \in \mathcal{L}^+(\mathcal{H}_A)$ $\text{Tr}_A(\rho_A) = 1$



Table: Composite systems

Classical Probability	Quantum Theory
Cartesian product $\Omega_{ZW} = \Omega_Z \times \Omega_W$	Tensor product $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$
Joint distribution $P(Z, W)$	Bipartite state $\rho_{AB}$
Marginal distribution $P(W) = \sum_{z \in \Omega_Z} P(Z = z, W)$	Reduced state $\rho_B = \text{Tr}_A(\rho_{AB})$
Conditional distribution $P(W Z) = \frac{P(Z, W)}{P(Z)}$	Conditional state $\rho_{B A} = ?$

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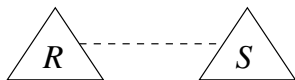


Figure: Classical correlations

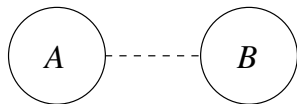


Figure: Quantum correlations

$$P(Z, W) = P(W|Z)P(Z)$$

$$\rho_{AB} = ?$$

## Definition

A **spatial quantum conditional state** of  $B$  given  $A$  is a positive operator  $\rho_{B|A}$  on  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$  that satisfies

$$\text{Tr}_B (\rho_{B|A}) = I_A.$$

c.f.  $P(W|Z)$  is a positive function on  $\Omega_{ZW} = \Omega_Z \times \Omega_W$  that satisfies

$$\sum_{w \in \Omega_W} P(W = w|Z) = 1.$$

$$(\rho_A, \rho_{B|A}) \quad \rightarrow \quad \rho_{AB} = (\sqrt{\rho_A} \otimes I_B) \rho_{B|A} (\sqrt{\rho_A} \otimes I_B)$$

$$\rho_{AB} \quad \rightarrow \quad \rho_A = \text{Tr}_B(\rho_{AB})$$

$$\rho_{B|A} = \left( \sqrt{\rho_A^{-1}} \otimes I_B \right) \rho_{AB} \left( \sqrt{\rho_A^{-1}} \otimes I_B \right)$$

$$(\rho_A, \rho_{B|A}) \quad \rightarrow \quad \rho_{AB} = (\sqrt{\rho_A} \otimes I_B) \rho_{B|A} (\sqrt{\rho_A} \otimes I_B)$$

$$\rho_{AB} \quad \rightarrow \quad \rho_A = \text{Tr}_B(\rho_{AB})$$

$$\rho_{B|A} = \left( \sqrt{\rho_A^{-1}} \otimes I_B \right) \rho_{AB} \left( \sqrt{\rho_A^{-1}} \otimes I_B \right)$$

Note:  $\rho_{B|A}$  defined from  $\rho_{AB}$  is a QCS on  $\text{supp}(\rho_A) \otimes \mathcal{H}_B$ .

**Table:** Comparison of relations between joints, conditionals and marginals

Classical Probability	Quantum Theory
$P(Z, W) = P(W Z)P(Z)$ $P(W Z) = \frac{P(Z, W)}{P(Z)}$	$\rho_{AB} = (\sqrt{\rho_A} \otimes I_B) \rho_{B A} (\sqrt{\rho_A} \otimes I_B)$ $\rho_{B A} =$ $\left( \sqrt{\rho_A^{-1}} \otimes I_B \right) \rho_{AB} \left( \sqrt{\rho_A^{-1}} \otimes I_B \right)$

- Drop implied identity operators, e.g.

- $I_A \otimes M_{BC} N_{AB} \otimes I_C \quad \rightarrow \quad M_{BC} N_{AB}$

- $M_A \otimes I_B = N_{AB} \quad \rightarrow \quad M_A = N_{AB}$

- Define non-associative “product”

- $M \star N = \sqrt{NM} \sqrt{N}$



**Table:** Comparison of relations between joints, conditionals and marginals

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$P(Z, W) = P(W Z)P(Z)$ $P(W Z) = \frac{P(Z, W)}{P(Z)}$	$\rho_{AB} = (\sqrt{\rho_A} \otimes I_B) \rho_{B A} (\sqrt{\rho_A} \otimes I_B)$ $\rho_{B A} =$ $\left( \sqrt{\rho_A^{-1}} \otimes I_B \right) \rho_{AB} \left( \sqrt{\rho_A^{-1}} \otimes I_B \right)$

**Table:** Comparison of relations between joints, conditionals and marginals

Classical Probability	Quantum Theory
$P(Z, W) = P(W Z)P(Z)$	$\rho_{AB} = \rho_{B A} \star \rho_A$
$P(W Z) = \frac{P(Z, W)}{P(Z)}$	$\rho_{B A} = \rho_{AB} \star \rho_A^{-1}$

- Given a classical variable  $Z$ , define a Hilbert space  $\mathcal{H}_Z$  with a preferred basis  $\{|1\rangle_Z, |2\rangle_Z, \dots, |d_Z\rangle_Z\}$  labeled by elements of  $\Omega_Z$ . Then,

$$\rho_Z = \sum_{z \in \Omega_Z} P(Z = z) |z\rangle \langle z|_Z$$

- Similarly,

$$\rho_{ZW} = \sum_{z \in \Omega_Z, w \in \Omega_W} P(Z = z, W = w) |zw\rangle \langle zw|_{ZW}$$

$$\rho_{W|Z} = \sum_{z \in \Omega_Z, w \in \Omega_W} P(W = w | Z = z) |zw\rangle \langle zw|_{ZW}$$

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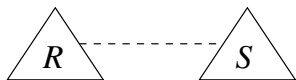


Figure: Classical correlations

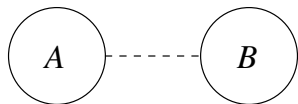


Figure: Quantum correlations

$$P(Z, W) = P(W|Z)P(Z)$$

$$\rho_{AB} = \rho_{B|A} \star \rho_A$$

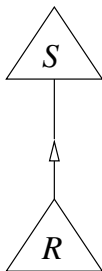


Figure: Classical stochastic map

$$\begin{aligned}
 P(W) &= \Gamma_{W|Z}(P(Z)) \\
 &= \sum_Z P(W|Z)P(Z)
 \end{aligned}$$

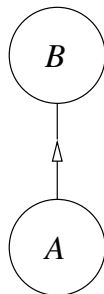


Figure: Quantum CPT map

$$\begin{aligned}
 \rho_B &= \mathcal{E}_{B|A}(\rho_A) \\
 &= \text{Tr}_A(\rho_{B|A} \star \rho_A) ?
 \end{aligned}$$

## Theorem (Jamiolkowski isomorphism)

Let  $\mathcal{E}_{B|A} : \mathcal{L}(\mathcal{H}_A) \rightarrow \mathcal{L}(\mathcal{H}_B)$  be a CPT map and define

$$\varrho_{B|A} = \mathcal{E}_{B|A'} \otimes \mathcal{I}_A \left( |\Phi^+\rangle \langle \Phi^+|_{A'A}^T \right),$$

where  $|\Phi^+\rangle_{A'A} = \sum_j |jj\rangle_{A'A}$ . Then,

$$\mathcal{E}_{B|A}(\rho_A) = \text{Tr}_A(\varrho_{B|A} \star \rho_A).$$

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- $\varrho_{B|A} = \rho_{B|A}^{T_A}$  for some spatial QCS  $\rho_{B|A}$ . Call such operators **temporal conditional states**.



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- $\varrho_{B|A} = \rho_{B|A}^{T_A}$  for some spatial QCS  $\rho_{B|A}$ . Call such operators **temporal conditional states**.
- Can also define **temporal joint state**:  $\varrho_{AB} = \varrho_{B|A} \star \rho_A$ .

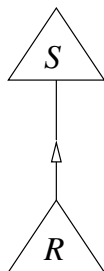


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 \end{aligned}$$

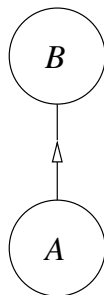


Figure: Quantum CPT map

$$\begin{aligned}
 \rho_B &= \mathcal{E}_{B|A}(\rho_A) \\
 &= \text{Tr}_A(\varrho_{B|A} \star \rho_A)
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- Composite of a quantum region and a classical random variable.
- Classical r.v.  $X$  has Hilbert space  $\mathcal{H}_X$  with preferred basis  $\{|1\rangle_X, |2\rangle_X, \dots, |d_X\rangle_X\}$ .
- Quantum region  $A$  has Hilbert space  $\mathcal{H}_A$ .
- Hybrid has Hilbert space  $\mathcal{H}_{XA} = \mathcal{H}_X \otimes \mathcal{H}_A$

- Composite of a quantum region and a classical random variable.
- Classical r.v.  $X$  has Hilbert space  $\mathcal{H}_X$  with preferred basis  $\{|1\rangle_X, |2\rangle_X, \dots, |d_X\rangle_X\}$ .
- Quantum region  $A$  has Hilbert space  $\mathcal{H}_A$ .
- Hybrid has Hilbert space  $\mathcal{H}_{XA} = \mathcal{H}_X \otimes \mathcal{H}_A$
- Operators on  $\mathcal{H}_{XA}$  restricted to be of the form

$$M_{XA} = \sum_{x \in \Omega_X} |x\rangle \langle x|_X \otimes M_{X=x,A}$$

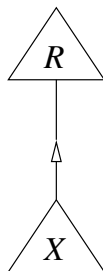


Figure: Classical preparation

$$P(Z) = \sum_X P(Z|X)P(X)$$

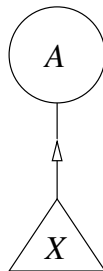


Figure: Quantum preparation

$$\rho_A = \sum_x P(X = x)\rho_A^{(x)}$$

$$\rho_A = \text{Tr}_X (\varrho_{A|X} \star \rho_X)$$

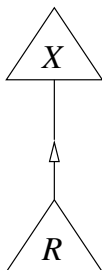


Figure: Noisy measurement

$$P(Y) = \sum_Z P(Y|Z)P(Z)$$

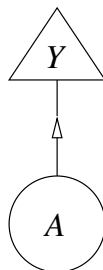


Figure: POVM measurement

$$P(Y = y) = \text{Tr}_A \left( E_A^{(y)} \rho_A \right)$$

$$\rho_Y = \text{Tr}_A \left( \rho_{Y|A} \star \rho_A \right)$$

## Definition (Quantum Instrument)

A **quantum instrument** is a set of CP-maps  $\mathcal{E}_{B|A}^{(y)}$  such that the operators  $E_A^{(y)}$  form a POVM, where

$$E_A^{(y)} = \left( \mathcal{E}_{B|A}^{(y)} \right)^\dagger (I_B).$$

- On obtaining  $y$  in a measurement of  $E_A^{(y)}$ :

$$\rho_A \rightarrow \rho_B^{(y)} = \frac{\mathcal{E}_{B|A}^{(y)}(\rho_A)}{\text{Tr}_A \left( E_A^{(y)} \rho_A \right)}$$

- The projection postulate is the special case:

$$\mathcal{E}_{B|A}^{(y)}(\rho_A) = \mathcal{I}_{B|A} \left( \Pi_A^{(y)} \rho_A \Pi_A^{(y)} \right)$$



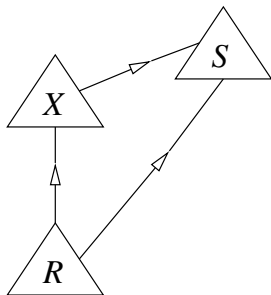


Figure: Classical instrument

$$P(Y, W) = \sum_Z P(Y, W|Z)P(Z)$$

$$P(Y|Z) = \sum_W P(Y, W|Z)$$

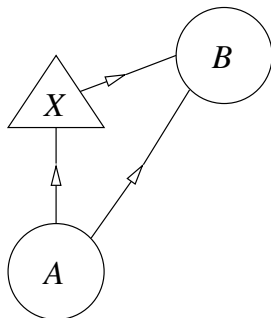


Figure: Quantum instrument

$$\rho_{YB} = \text{Tr}_A (\varrho_{YB|A} \star \rho_A)$$

$$\varrho_{Y|A} = \text{Tr}_B (\varrho_{YB|A})$$

Dynamics	$\rho_B = \mathcal{E}_{B A}(\rho_A)$	$\rho_B = \text{Tr}_A(\varrho_{B A} \star \rho_A)$
Preparation	$\rho_A = \sum_x P(X=x)\rho_A^{(x)}$	$\rho_A = \text{Tr}_X(\varrho_{A X} \star \rho_X)$
Measurement	$P(Y=y) = \text{Tr}_A(E^{(y)}\rho_A)$	$\rho_Y = \text{Tr}_A(\varrho_{Y A} \star \rho_A)$
Update	$P(Y=y)\rho_B^{(y)} = \mathcal{E}_{B A}^{(y)}(\rho_A)$	$\rho_{YB} = \text{Tr}_A(\varrho_{YB A} \star \rho_A)$

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- Two expressions for joint probabilities:

$$\begin{aligned}P(Z, W) &= P(W|Z)P(Z) \\ &= P(Z|W)P(W)\end{aligned}$$

- Bayes' rule:

$$P(Z|W) = \frac{P(W|Z)P(Z)}{P(W)}$$

- Alternative form of Bayes' rule:

$$P(Z|W) = \frac{P(W|Z)P(Z)}{\sum_Z P(W|Z)P(Z)}$$

- Two expressions for bipartite states:

$$\begin{aligned}\rho_{AB} &= \rho_{B|A} \star \rho_A \\ &= \rho_{A|B} \star \rho_B\end{aligned}$$

- Bayes' rule:

$$\rho_{A|B} = \rho_{B|A} \star \left( \rho_A \otimes \rho_B^{-1} \right)$$

- Alternative form of Bayes' rule

$$\rho_{A|B} = \rho_{B|A} \star \left( \rho_A \otimes \text{Tr}_A (\rho_{B|A} \star \rho_A)^{-1} \right)$$

- Given an input state  $\rho_A$  and a temporal QCS  $\varrho_{B|A}$ , define

$$\varrho_{A|B} = \varrho_{B|A} \star (\rho_A \otimes \text{Tr}_B (\varrho_{B|A} \star \rho_A))$$

- Temporal joint state now has two decompositions:

$$\varrho_{AB} = \varrho_{B|A} \star \rho_A = \varrho_{A|B} \star \rho_B$$

- A temporal hybrid joint state can be written two ways:

$$\varrho_{XA} = \varrho_{A|X} \star \rho_X = \varrho_{X|A} \star \rho_A$$

- The two representations are connected via Bayes' rule:

$$\varrho_{X|A} = \varrho_{A|X} \star \left( \rho_X \otimes \text{Tr}_X (\varrho_{A|X} \star \rho_X)^{-1} \right)$$

$$\varrho_{A|X} = \varrho_{X|A} \star \left( \text{Tr}_A (\varrho_{X|A} \star \rho_A)^{-1} \otimes \rho_A \right)$$

$$\varrho_{X=x|A} = \frac{P(X=x)\varrho_{A|X=x}}{\sum_{x' \in \Omega_X} P(X=x')\varrho_{A|X=x'}}$$

$$\varrho_{A|X=x} = \frac{\sqrt{\rho_A}\varrho_{X=x|A}\sqrt{\rho_A}}{\text{Tr}_A (\varrho_{X=x|A}\rho_A)}$$

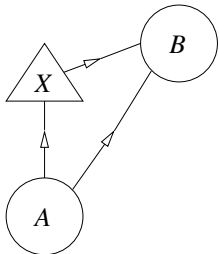
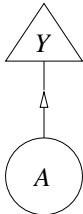
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- Classically, upon learning  $X = x$ :

$$P(Z) \rightarrow P(Z|X = x)$$

- Quantumly:  $\rho_A \rightarrow \varrho_{A|X=x}$ ?

Projection Postulate	Bayesian Conditioning
$\rho_A \rightarrow \frac{\mathcal{I}_{B A} \left( \sqrt{E_A^{(y)}} \rho_A \sqrt{E_A^{(y)}} \right)}{\text{Tr}_A \left( E_A^{(y)} \rho_A \right)}$  <pre> graph BT     A((A)) --&gt; X(△X)     A((A)) --&gt; B((B))     B((B)) --&gt; X(△X)     </pre>	$\rho_A \rightarrow \frac{\sqrt{\rho_A} E_A^{(y)} \sqrt{\rho_A}}{\text{Tr}_A \left( E_A^{(y)} \rho_A \right)}$  <pre> graph BT     A((A)) --&gt; Y(△Y)     </pre>

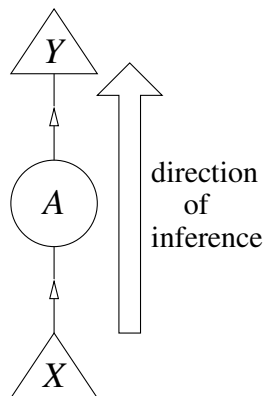


Figure: Prep. & meas. experiment

- Joint probability:  

$$\varrho_{XY} = \text{Tr}_A (\varrho_{Y|A} \star (\varrho_{A|X} \star \rho_X))$$
- Marginal for Y:  

$$\rho_Y = \text{Tr}_A (\varrho_{Y|A} \star \rho_A)$$
- Conditional probabilities:  

$$\varrho_{Y|X} = \text{Tr}_A (\varrho_{Y|A} \star \varrho_{A|X})$$
- Bayesian update:  

$$\rho_A \rightarrow \varrho_{A|X=x}$$

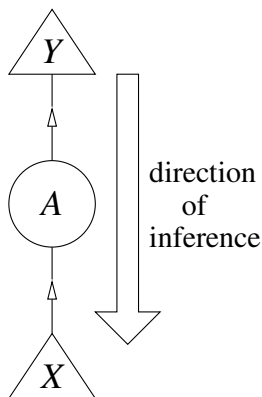


Figure: Prep. & meas. experiment

- Apply Bayes' rule to  $\varrho_{A|X}$  and  $\varrho_{Y|A}$ :

$$\rho_{XY} = \text{Tr}_A (\varrho_{X|A} \star (\varrho_{A|Y} \star \rho_Y))$$

- Marginal for  $X$ :

$$\rho_X = \text{Tr}_A (\varrho_{X|A} \star \rho_A)$$

- Conditional probabilities:

$$\varrho_{X|Y} = \text{Tr}_A (\varrho_{X|A} \star \varrho_{A|Y})$$

- Bayesian update:

$$\rho_A \rightarrow \varrho_{A|Y=y}$$

- c.f. Barnett, Pegg & Jeffers, J. Mod. Opt. 47:1779 (2000).

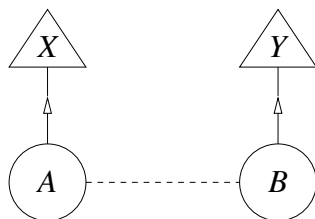


Figure: Bipartite experiment

- Joint probability:  $\rho_{XY} = \text{Tr}_{AB} ((\varrho_{X|A} \otimes \varrho_{Y|B}) \star \rho_{AB})$
- $B$  can be factored out:  $\rho_{XY} = \text{Tr}_A (\varrho_{Y|A} \star (\varrho_{A|X} \star \rho_X))$
- where  $\varrho_{Y|A} = \text{Tr}_B (\varrho_{Y|B} \rho_{B|A})$

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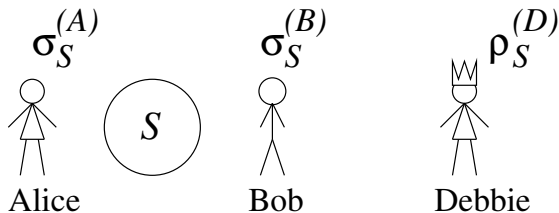


Figure: Initial State Assignments

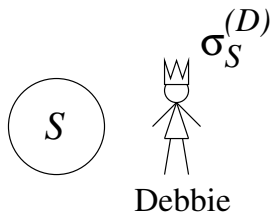


Figure: Final State Assignment

$$\sigma_S^{(D)} = f(\sigma_S^{(A)}, \sigma_S^{(B)}, \rho_S^{(D)})$$



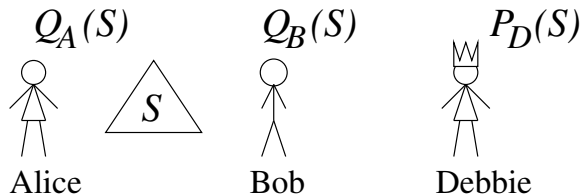


Figure: Initial State Assignments

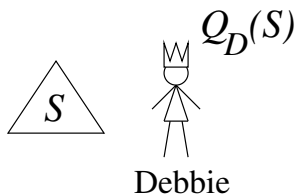


Figure: Final State Assignment

$$Q_D(S) = f(Q_A(S), Q_B(S), P_D(S))$$

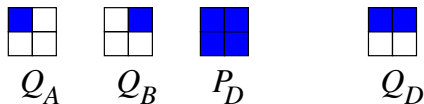


Figure: Combining Incompatible Assignments

- Linear pool:

$$Q_D(S) = w_A Q_A(S) + w_B Q_B(S) + w_D P_D(S)$$

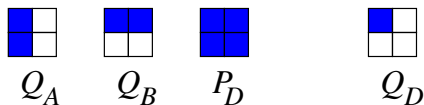


Figure: Combining Independent Evidence

- Multiplicative (log-linear) pool:

$$Q_D(S) \propto \frac{Q_A(S)Q_B(S)}{P_D(S)}$$

- Bayesian inference says that:

$$\begin{aligned} Q_D(S) &= P_D(S | R_A = Q_A(S), R_B = Q_B(S)) \\ &= \frac{P_D(R_A = Q_A(S), R_B = Q_B(S) | S) P_D(S)}{\sum_S P_D(R_A = Q_A(S), R_B = Q_B(S) | S) P_D(S)} \end{aligned}$$

- Similarly:

$$\begin{aligned} \sigma_S^{(D)} &= \rho_{S | R_A = \sigma_S^{(A)}, R_B = \sigma_S^{(B)}}^{(D)} \\ &= \rho_{R_A = \sigma_S^{(A)}, R_B = \sigma_S^{(B)} | S}^{(D)} \star \left( \rho_S^{(D)} \otimes \text{Tr}_S \left( \rho_{R_A = \sigma_S^{(A)}, R_B = \sigma_S^{(B)} | S}^{(D)} \star \rho_S^{(D)} \right) \right) \end{aligned}$$

- Quantum Chain Rule:

$$\rho_{ABC} = \rho_{C|AB} \star (\rho_{B|A} \star \rho_A)$$

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## Definition

If  $\rho_{C|AB} = \rho_{C|B}$  then  $C$  is **conditionally independent** of  $A$  given  $B$ .

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## Definition

If  $\rho_{C|AB} = \rho_{C|B}$  then  $C$  is **conditionally independent** of  $A$  given  $B$ .

## Theorem

*The following conditions are equivalent:*

- $\rho_{C|AB} = \rho_{C|B}$
- $\rho_{A|BC} = \rho_{A|B}$
- $I(A : C|B) = 0$ .

*Further, conditional independence implies that*

- $\rho_{AC|B} = \rho_{A|B}\rho_{C|B}$ .



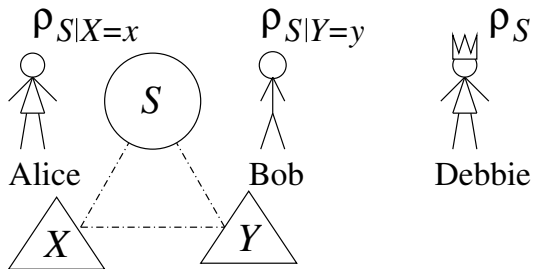


Figure: The Case of Shared Priors

## Theorem

If  $X$  and  $Y$  are conditionally independent given  $S$  then

$$\rho_{S|R_A=\rho_{S|X=x}, R_B=\rho_{S|Y=y}}^{(D)} \propto \rho_{S|X=x} \rho_S^{-1} \rho_{S|Y=y}$$

### Theorem (Stronger Version)

*If the minimal sufficient statistics for  $X$  and  $Y$  with respect to  $S$  are conditionally independent given  $S$  then*

$$\rho_{S|R_A=\rho_{S|X=x}, R_B=\rho_{S|Y=y}}^{(D)} \propto \rho_{S|X=x} \rho_S^{-1} \rho_{S|Y=y}$$

- 1 Introduction
- 2 Spatial Quantum Conditional States
- 3 Temporal Quantum Conditional States
- 4 Preparations and Measurements
- 5 Quantum Bayes' rule
- 6 Bayesian Conditioning
- 7 Combining Expert Advice
- 8 Further results and open questions**

## Forthcoming paper(s) with R. W. Spekkens also include:

- Quantum sufficient statistics
- Quantum state compatibility
- Quantum pooling

## Earlier papers with related ideas:

- M. Asorey et. al., *Open.Syst.Info.Dyn.* 12:319–329 (2006).
- M. S. Leifer, *Phys. Rev. A* 74:042310 (2006).
- M. S. Leifer, *AIP Conference Proceedings* 889:172–186 (2007).
- M. S. Leifer & D. Poulin, *Ann. Phys.* 323:1899 (2008).

What is the meaning of fully quantum Bayesian conditioning?

$$\rho_B \rightarrow \rho_{B|A} = \rho_{A|B} \star \left( \text{Tr}_B (\rho_{A|B} \star \rho_B)^{-1} \otimes \rho_B \right)$$

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- University College London